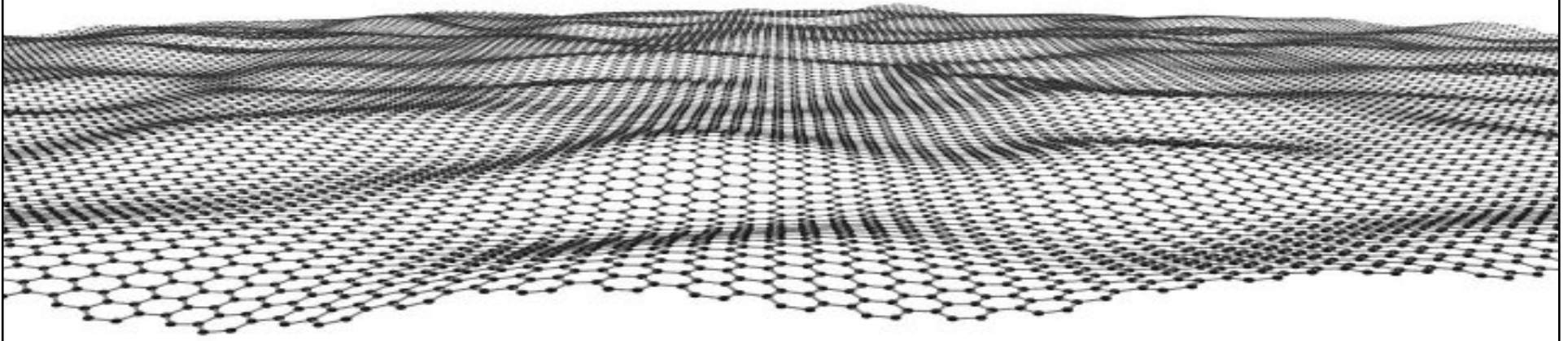


Theory-Seminar: Graphene

Phonons, Elasticity and Crumpling of Graphene





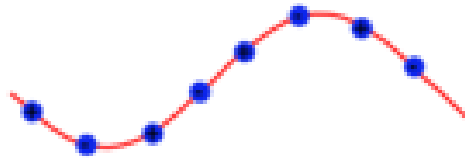
**GRAPHENE -
TO BE OR NOT TO BE
... IN 2D?**

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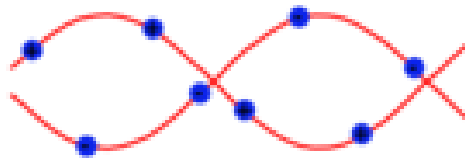
- I. Introduction: Phonons in 3D and 2D
- II. Phonons in graphene
- III. Dispersion relation of free floating graphene-sheets
- IV. Dispersion relation of graphene in presence of tension or a substrate
- V. Number of phonons and structure
- VI. Approximations
- VII. Summary

Phonons in crystal lattices

Akustische Mode



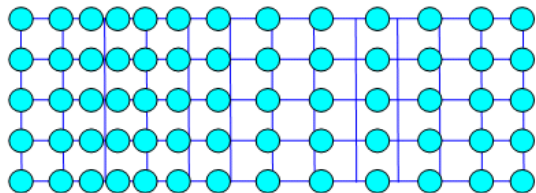
Optische Mode



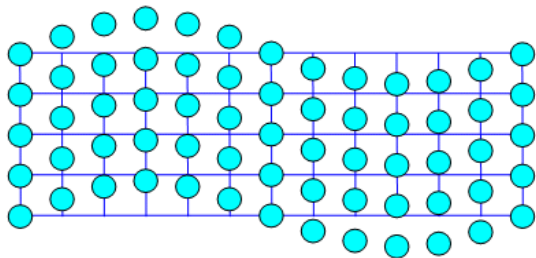
Branches of a **lattice base with r atoms**
in **D dimensions:**

D acoustic branches

$D \cdot r - D$ optical branches

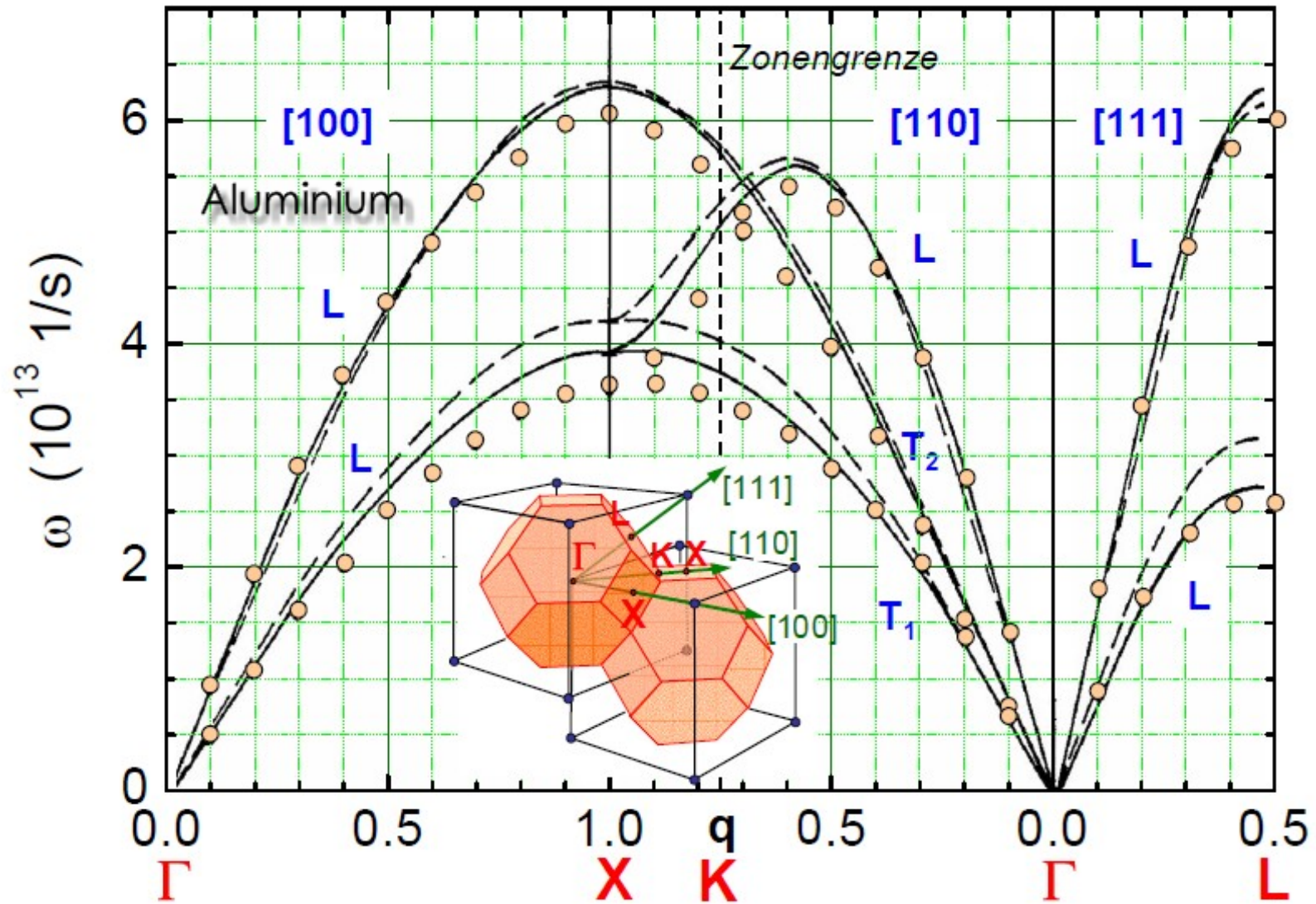


Longitudinal mode

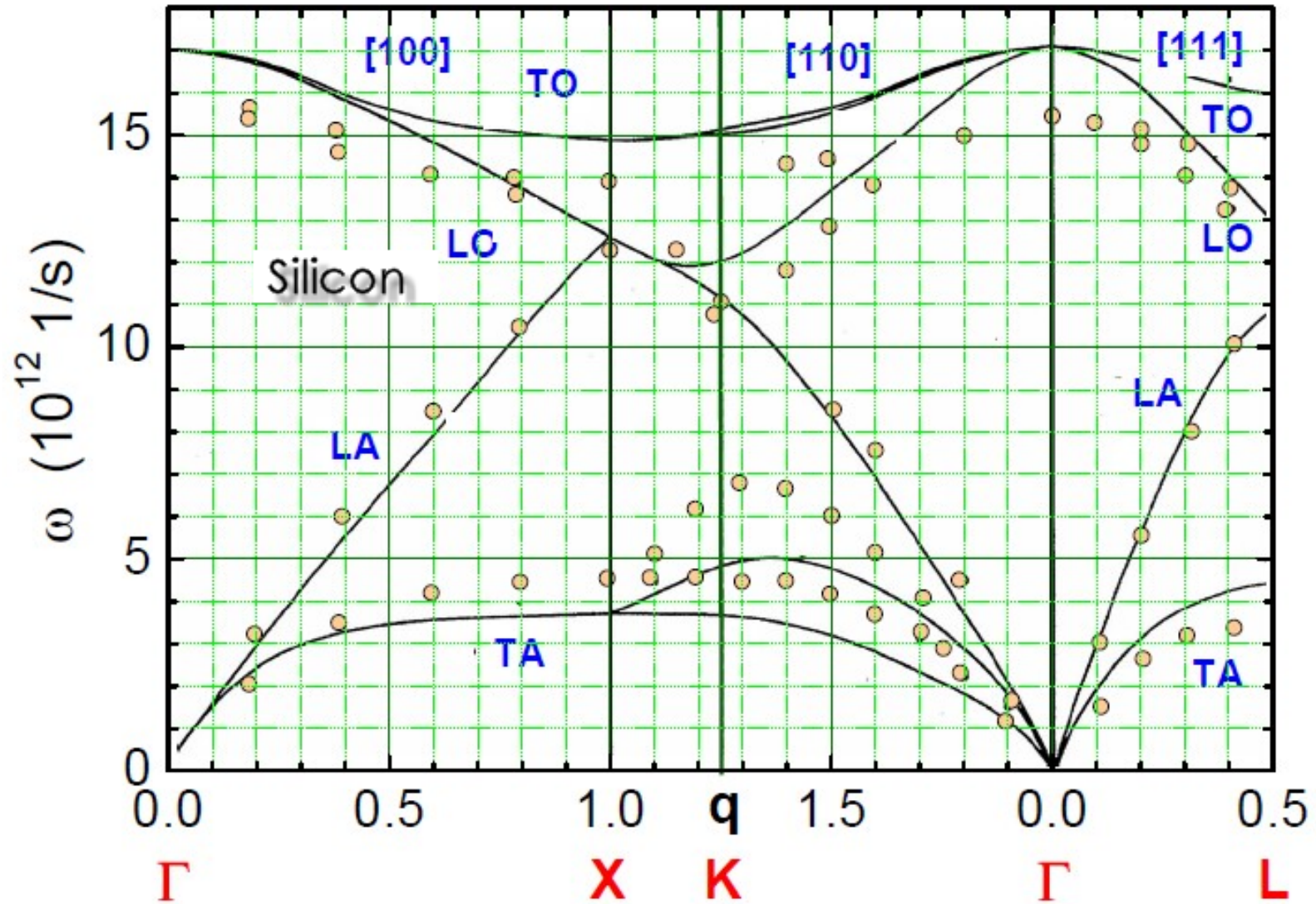


transversal mode

Phonons in 3D (one atom in base)



Phonons in 3D (two atom in base)



Phonons in 3D (four atom in base)

3 acoustic branches

$3 \cdot 4 - 3 = 9$ optical branches

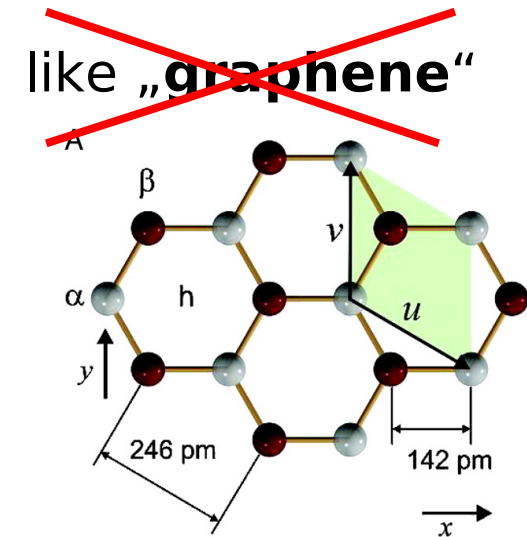
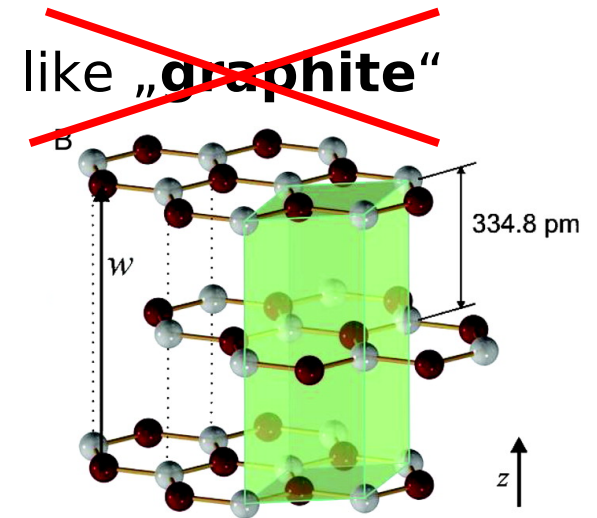
12 branches

Phonons in 2D (two atom in base)

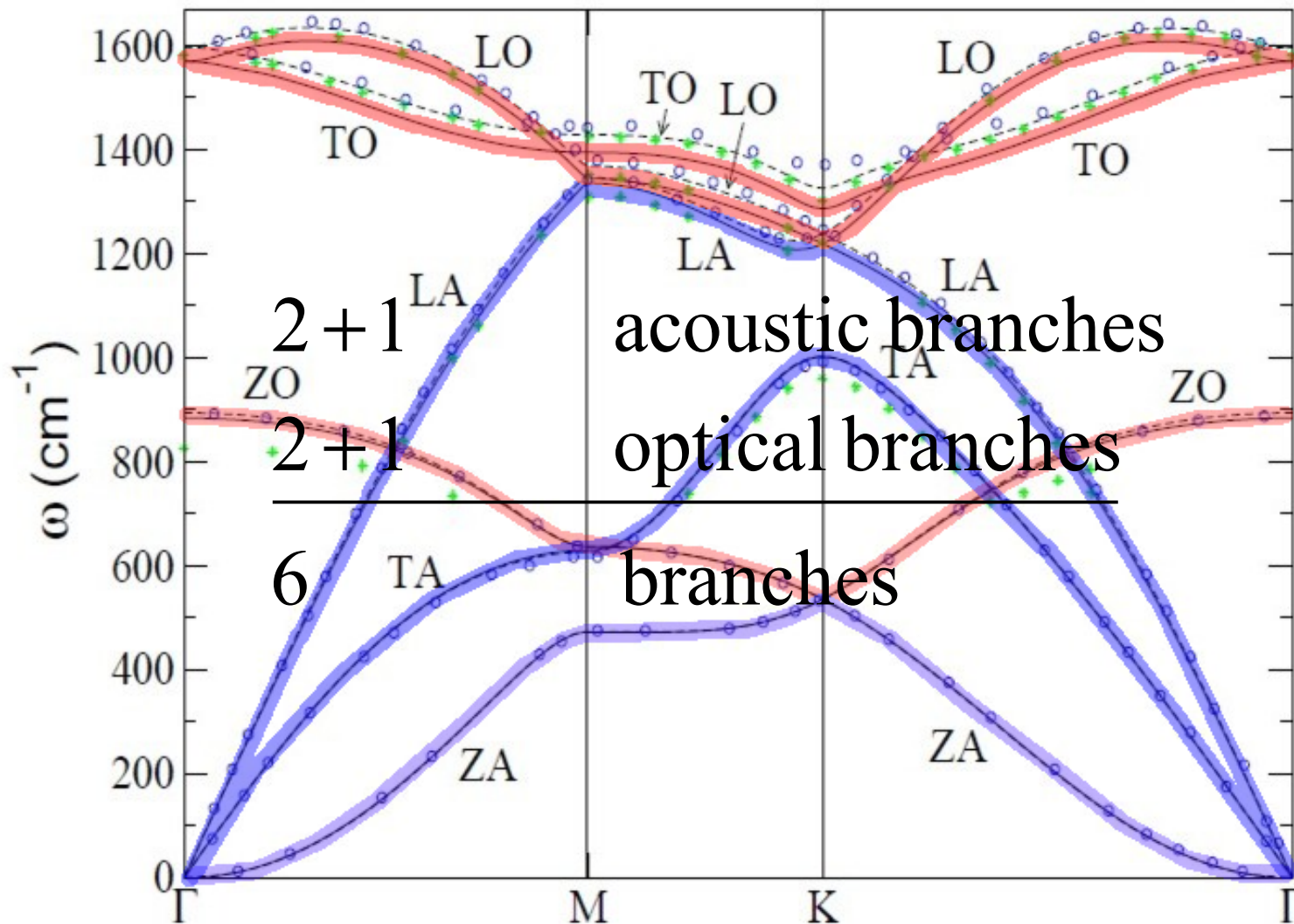
2 acoustic branches

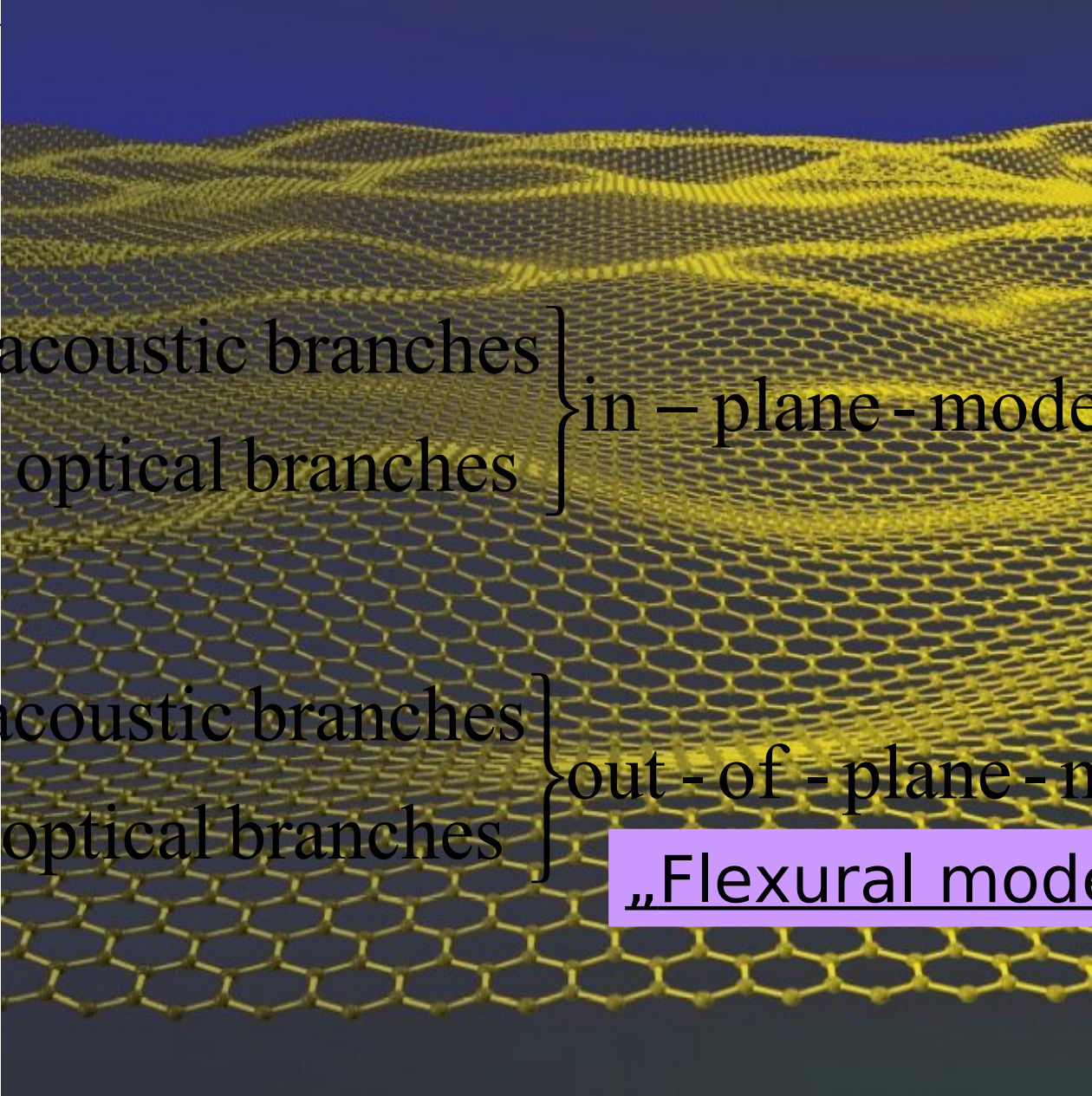
$2 \cdot 2 - 2 = 2$ optical branches

4 branches



Phonons in Graphene





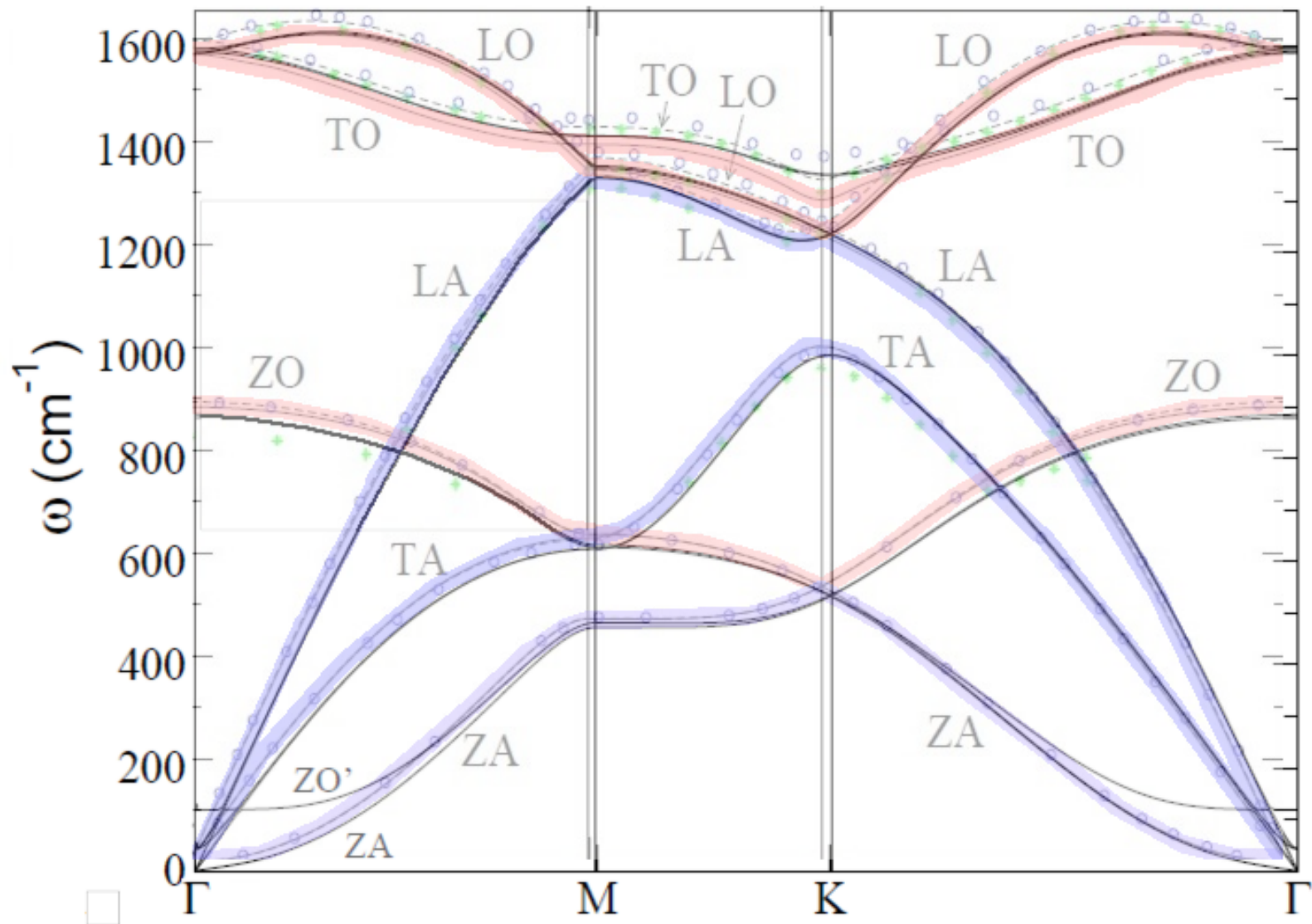
$\left\{ \begin{array}{l} 2 \text{ acoustic branches} \\ 2 \text{ optical branches} \end{array} \right\}$ in - plane - modes

+

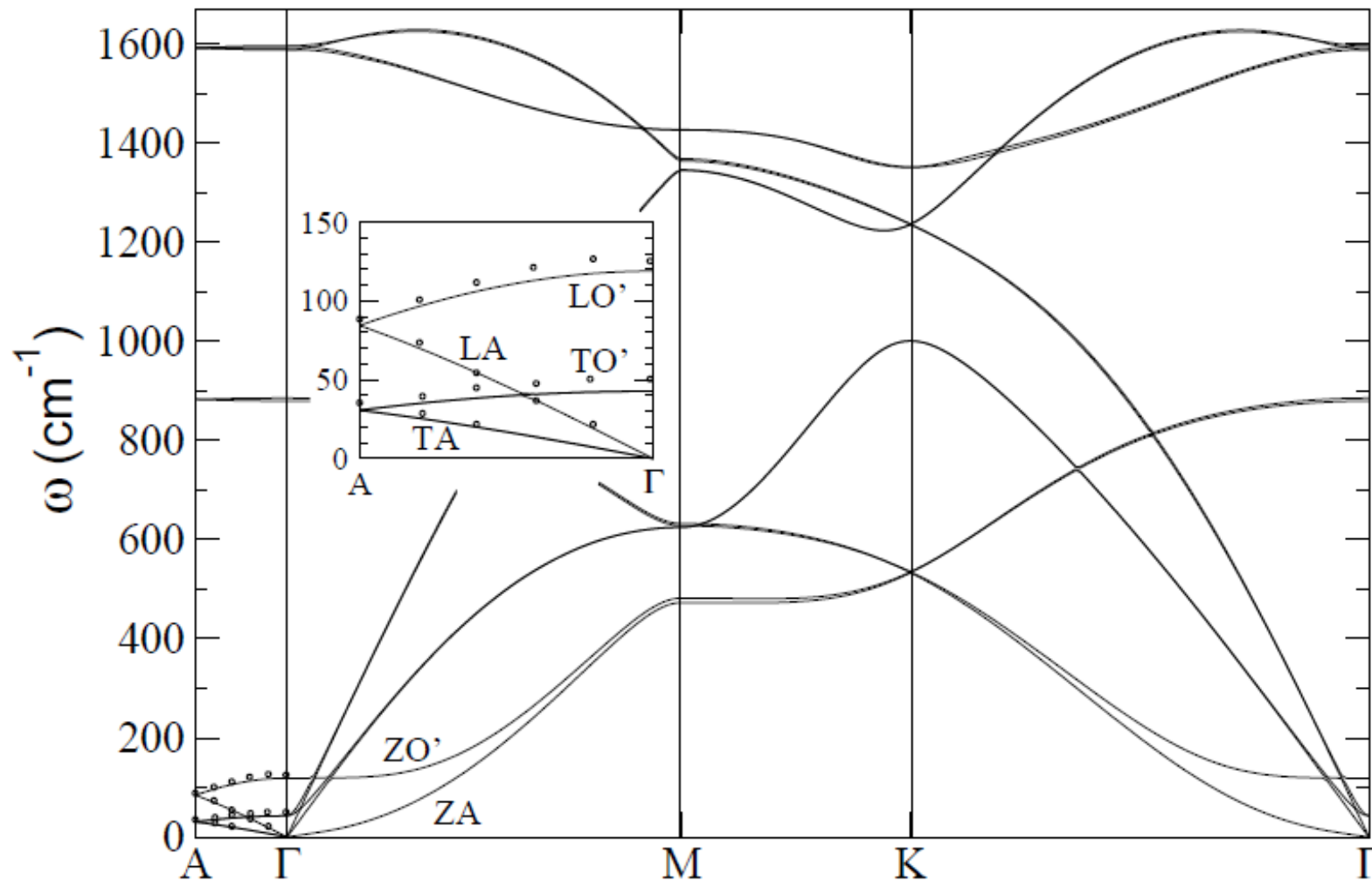
$\left\{ \begin{array}{l} 1 \text{ acoustic branches} \\ 1 \text{ optical branches} \end{array} \right\}$ out - of - plane - modes

„Flexural modes“

Phonons in Graphite



Phonons in Graphite



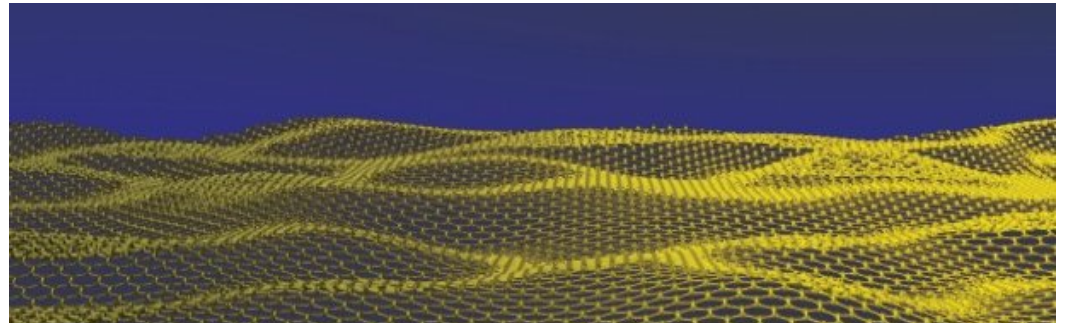
Dispersion Relation of Phonons in Graphene

Procedure:

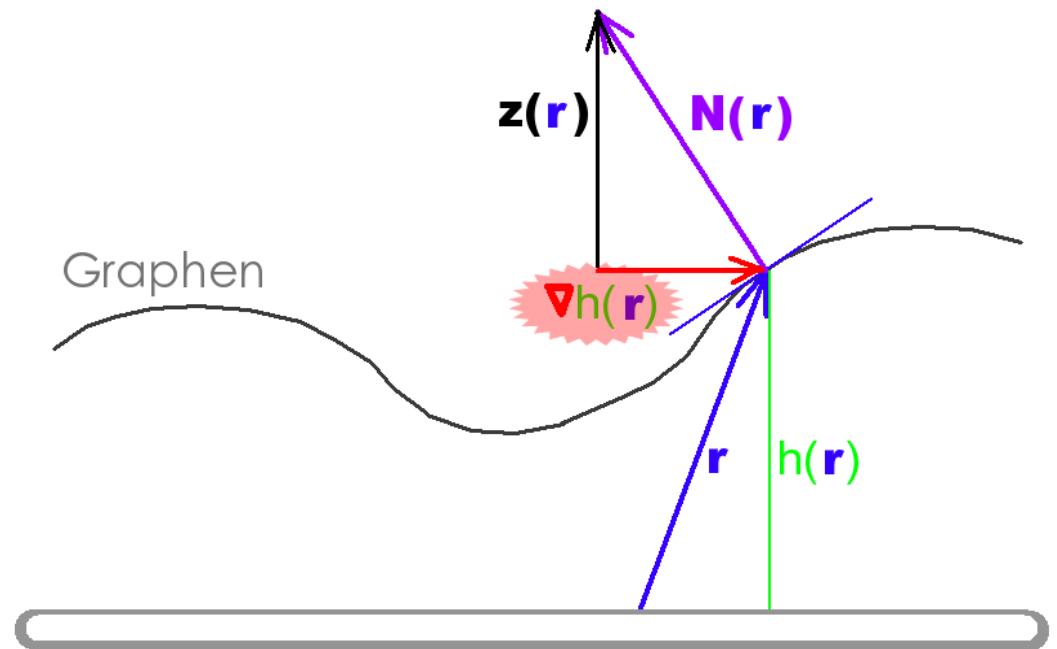
- Create a **potential** for an oscillation in free floating graphene
 - Put this potential into a **Hamilton-Operator**
 - Solve the **equations of motion**
-
- Adapt the dispersion relation to **tension and substrate**
-
- Calculate the **number of phonons**
-
- Think about all the **approximations** you have done!

1. Potential for an oscillation

How to describe the elongation?



$$N = \frac{z - \nabla h}{\sqrt{1 + (\nabla h)^2}}$$



1. Potential for an oscillation

Integration over microscopic springs:

$$E_0 = \frac{\mathcal{K}}{2} \int d^2 r (\nabla N)^2$$

For small bendings $(\nabla h)^2 \ll 1$ (**harmonic approximation**) we obtain

$$E_0 \approx \frac{\mathcal{K}}{2} \int d^2 r (\nabla^2 h)^2$$

Fourier-Transformation

$$h(\mathbf{r}) = \frac{1}{2\pi} \int d^2 k e^{i\mathbf{k}\cdot\mathbf{r}} h_{\mathbf{k}}$$

2. Hamilton-Operator

Define a momentum operator and put kinetic and the potential-energy term into the Hamilton-Operator

$$H = \sum_k \left\{ \frac{1}{2\sigma} P_{-k} P_k + \frac{\kappa k^4}{2} h_{-k} h_k \right\}$$

Compare with Hamiltonian of classical spring problem

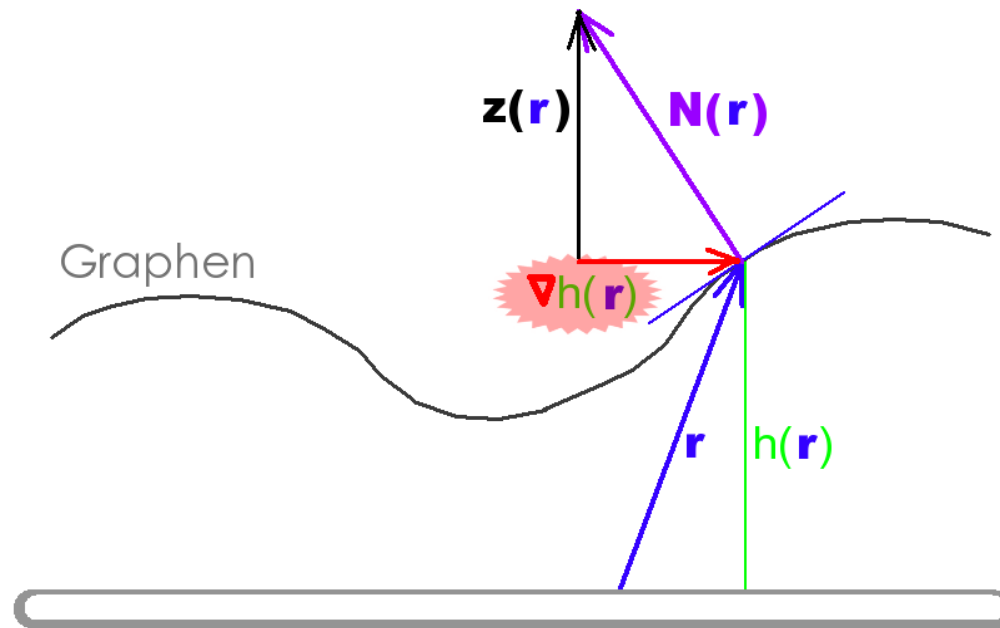
$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \quad \Rightarrow \quad \sigma \omega^2 = \kappa k^4$$

$$\Rightarrow \omega_{flex}(k) = \sqrt{\frac{\kappa}{\sigma}} \cdot k^2 \quad \text{Dispersion relation for free floating graphene}$$

3. Graphene under tension

Integration over microscopic springs:

$$E_T = \frac{\gamma}{2} \int d^2r (\nabla h)^2$$



3. Graphene under tension

Same calculation for graphene under tension ending up in a different Hamilton-Operator

$$H = \sum_k \left\{ \frac{1}{2\sigma} P_{-k} P_k + \frac{\kappa k^4}{2} h_{-k} h_k + \frac{\gamma k^2}{2} h_{-k} h_k \right\}$$

Compare with Hamiltonian of classical spring problem

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

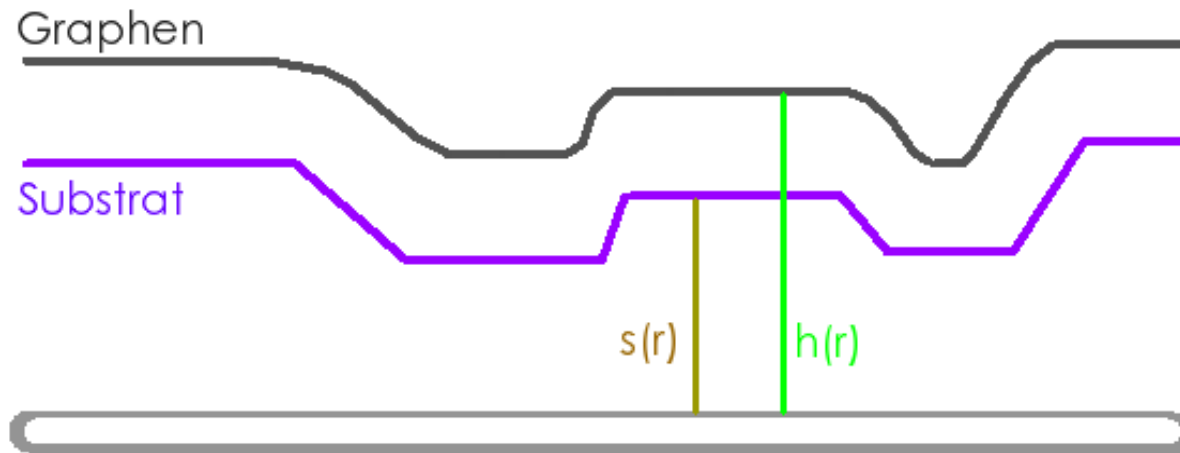
$$\Rightarrow \omega(k) = k \sqrt{\frac{\kappa}{\sigma} k^2 + \frac{\gamma}{\sigma}}$$

Dispersion relation for graphene under tension

4. Graphene on a substrate

Graphene in presence of a substrate

$$E_s \approx \frac{v}{2} \int d^2r (s(r) - h(r))^2$$



5. Graphene as a Graphite layer

Is there any relation between the stiffness of graphene and graphite?

Weak Van-der-Waals forces between graphite-layers,
thus a relation is possible

Lets measure the resonance frequency of a
microscopic graphite sample

$$v(k) = \sqrt{\frac{Y}{\sigma}} t k^2$$

Compare $v(k)$ to $\omega_{flex}(k) = \sqrt{\frac{\kappa}{\sigma}} \cdot k^2$


$$\kappa = Yt^3$$

This value fits to the measurements with graphene

6. Number of Phonons

a) Free floating graphene

$$N_{ph} = \int \frac{d^2k}{(2\pi)^2} \cdot n_k$$
$$= \int_0^\infty \frac{dk}{2\pi} \cdot \frac{k}{\exp(\beta \sqrt{\kappa/\sigma} \cdot k^2) - 1}$$

$$n_k = \frac{1}{\exp(\beta(E)) - 1}$$

$$\omega_{flex}(k) = \sqrt{\frac{\kappa}{\sigma}} \cdot k^2$$

Integral diverges in the infrared!!

That means graphene could not exist!!

Way out: Set the lower limit $2\pi/L$

6. Number of Phonons

Now the integral converges

$$N_{ph} = \frac{\pi}{L_T} \ln \left(\frac{1}{1 - \exp(-L_T / L)} \right)$$

With the thermal wavelength $L_T = \frac{\pi}{\sqrt{k_B T}} \left(\frac{\kappa}{\sigma} \right)^{1/2}$

6. Number of Phonons

$$N_{ph} = \frac{\pi}{L_T^2} \ln \left(\frac{1}{1 - \exp(-L_T^2 / L^2)} \right)$$

Have a closer look at $\left(\frac{L_T}{L} \right)$ Keep in mind: $L_T \propto \frac{1}{\sqrt{T}}$

1. $L \gg L_T$

**Number of phonons diverges
Logarithmically with L**

**System cannot be structurally ordered
at any finite temperature**

6. Number of Phonons

$$N_{ph} = \frac{\pi}{L_T^2} \ln \left(\frac{1}{1 - \exp(-L_T^2/L^2)} \right)$$

Have a closer look at $\left(\frac{L_T}{L} \right)$ Keep in mind: $L_T \propto \frac{1}{\sqrt{T}}$

2. $L \ll L_T$

Number of phonons goes to 0

**System can be flat
if temperature is low enough**

6. Number of Phonons

What means „low enough temperature“?

1. T = 300K

One obtain $L_T \approx 0,1nm \Rightarrow L_T \ll L$

2. T = 0,01K

One obtain $L_T \approx 100nm \Rightarrow L_T \ll L$

The statement **System cannot be structurally ordered** won!
At any finite temperature

6. Number of Phonons

b) How does graphene look like **under tension**?

$$N_{ph} = \int \frac{d^2k}{(2\pi)^2} \cdot n_k$$

$$n_k = \frac{1}{\exp(\beta(E)) - 1}$$

Integral converges in the infrared.

That means it is independent from the size,

Just depends on the temperature.

Under these circumstances it is easier to find conditions, where

a graphene-sheet could be flat at finit temperatures.

$$\omega(k) = k \sqrt{\frac{\kappa}{\sigma} k^2 + \frac{\gamma}{\sigma}}$$

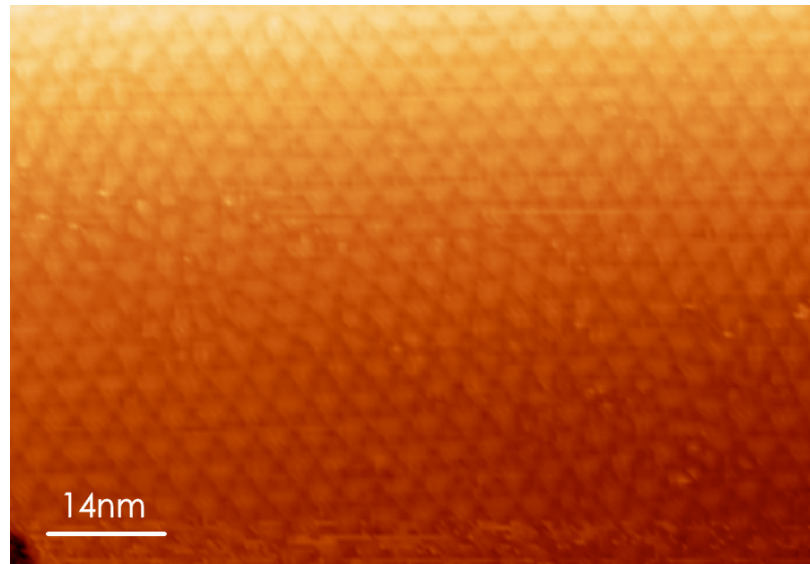
6. Number of Phonons

c) How does graphene look like on a substrate?

Substrate has a certain surface, described by $s(r)$.

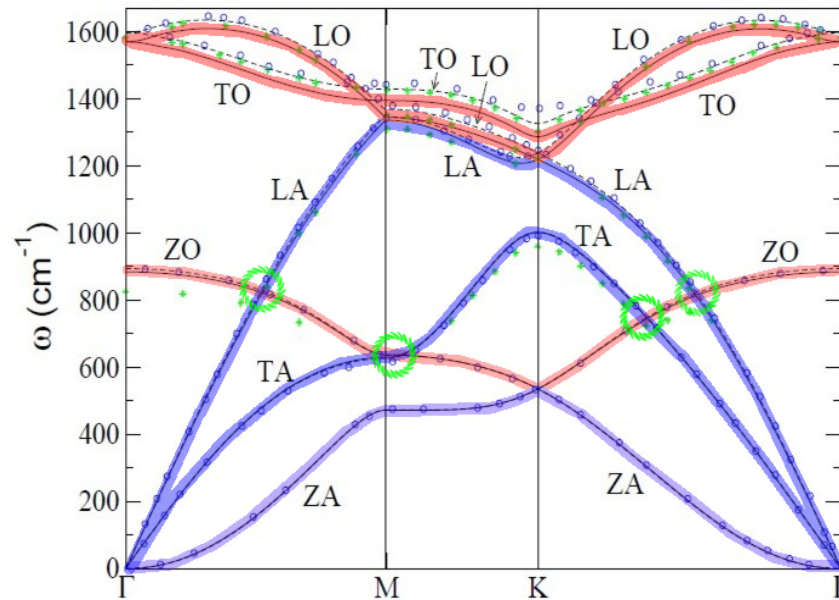
Graphene tries to follow the substrate surface to minimize potential energy.

Obviously, the graphene-sheet could have a static rippling.



7. Approximations and reality

1. We assumed that the **graphene-lattice is perfect**
BUT: there are topological defect: disclination, dislocation
2. We assumed **small bending deformations**
BUT: What's small?
3. We assumed that there is **no coupling between the modes**
BUT: Coupling exists



7. Approximations and reality

All these things lead to **non-linear effects**
and to such a **renormalization of the bending rigidity**

that a flat graphene-phase is possible at low temperatures

Finally there can be **electron-phonon-interactions** between a
metal gate and the flexural modes, which leads to a

damping of the flexural modes

Summary

- **Graphene does exist** - in theorie AND reality
- Graphene is a quasi-2D-system and has **flexural modes**
- One can get the dispersion relation by solving the linear problem of a **harmonic oscillator**
- There are different **dispersion relations** for different circumstances (**k^2 - free floating, linear in k - under tension**)
- **Number of phonons** tell us about the **structural order** of graphene
- **Free floating** graphene **cannot be flat**, but **not free floating** graphene **can be flat** under certain condition

Literature

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A.K. Geim, A.M.	Graphene - exploring carbon flatland [BibTeX]	2007	Physics Today Vol. August 2007, pp. 35-41	article
A.K. Geim, K.N.	The rise of graphene [BibTeX]	2007	Nature Materials Vol. 6, pp. 183-191	article
Katsnelson, M.	Graphene: carbon in two dimensions [BibTeX]	2007	Materials today Vol. 10, no. 1-2, pp. 20-27	article
R. Gross, A.M.	Festkörperphysik - Vorlesungsskript zur Vorlesung im WS 2004/2005 [BibTeX]	2004		booklet

End

